

Resit
First part (T1-T3) of
Thermal Physics
2020-2021

Saturday April 10, 2021
9.00-12.00
online

PROBLEM 1

Score: $a+b+c+d+e=6+6+6+6+6=30$

A container with a volume of 1.0 dm^3 contains $2.0 \cdot 10^{23}$ He atoms. The temperature of the He gas is 54.4 K . You can assume that the He behaves like an ideal gas.

Given: $M_{\text{He}} = 4.0 \text{ g mol}^{-1}$, $M_{\text{N}_2} = 28.0 \text{ g mol}^{-1}$, $M_{\text{O}_2} = 32.0 \text{ g mol}^{-1}$.

- a) Determine the gas pressure.
- b) What is the root mean square (RMS) speed of the molecules at this temperature?
- c) What is the internal energy of the gas in the container?
- d) Now you replace every He atom by a N_2 molecule. How do temperature, RMS speed and internal energy of the gas change? Explain.
- e) Draw a sketch of the Maxwell distribution. Roughly approximate the fraction of O_2 molecules that have molecular speeds in the range between 1500 m/s and 1600 m/s when the temperature is 300 K .

PROBLEM 2

Score: $a+b+c+d+e=6+6+6+6+6=30$

Consider a hypothetical engine which connects two (infinite) reservoirs and undergoes the following *reversible* processes:

- i. isothermal expansion at temperature T_h
- ii. adiabatic expansion to temperature T_c
- iii. isothermal compression at temperature T_c
- iv. adiabatic compression to the initial state

In every cycle, a quantity of heat Q_{in} flows from reservoir A into the engine and a quantity of heat Q_{out} flows from the engine into the reservoir B.

- a) Sketch the thermodynamic cycle in the usual (p,V) -diagram. Clearly indicate the direction of the cycle. In which steps does the heat flow and in which steps is the engine doing work?
- b) Sketch the thermodynamic cycle in a (slightly less usual) (T,S) -diagram. Indicate where the entropy of the engine changes and how these changes relate to the entropies of reservoirs A and B.
- c) Use entropy arguments to show, that the efficiency of the cycle equals $\eta = 1 - T_c/T_h$.
- d) The coal power plant in Eemshaven produces about 12 TWh of electricity per year from about 3 million tons of coal (thermal energy content of coal: 9 kWh kg^{-1}). Determine the efficiency of the power plant and estimate T_c and T_h assuming the power plant is an idealized Carnot engine. Is the result for T_h realistic? Estimate the efficiency of a primitive steam engine (use realistic numbers) and compare.
- e) Until now we have assumed infinite reservoirs A and B. Now assume both reservoirs are finite (and have the same heat capacity). Discuss the effect on the thermodynamic cycle and on the reservoir temperatures.

PROBLEM 3

Score: $a+b+c = 6+6+6=18$

The thermal diffusion equation for a sphere can be written as

$$\frac{\partial T}{\partial t} = D \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial T}{\partial r} \right)$$

if there is no dependence on the angular coordinates. In this equation r is the radial coordinate, D is the diffusion coefficient and T is the temperature.

- a) Give a general solution for the steady state case.
- b) Give the solution for the boundary condition of a spherical animal of radius r_0 , with body temperature T_{body} for $r = r_0$ and a temperature of the outside medium $T_{outside}$ for $r \rightarrow \infty$.
- c) How much heat does the animal loose to the medium per second? The thermal conductivity of the medium is κ .

PROBLEM 4

Score: $a+b=6+6=12$

Typical turbomolecular pumps can generate a vacuum of about $1 \cdot 10^{-10}$ bar. Assume the pump to be at $T = 25^\circ C$ and working on a gas primarily consisting of N_2 molecules.

- a) Calculate the mean free path of the molecules and the collision frequency.
- b) Do the molecules that pass the turbomolecular pump have the same Maxwell Boltzmann distribution of speeds that the gas in the vacuum chamber has? (Hint: assume that the pump has a diameter of 10 cm. Use the result for the mean free path to argue.) If the answer is no, what speed distribution do these molecules have?

Collision diameter N_2 : 395 pm.

Solutions

PROBLEM 1

a)

$$pV = Nk_B T$$

$$p = \frac{Nk_B T}{V} = \frac{2 \times 10^{23} \times 1.38 \times 10^{-23} \text{ J/K} \times 54.4 \text{ K}}{10^{-3} \text{ m}^3} \approx 150 \text{ kPa}$$

b)

$$v_{rms} = \sqrt{\langle v^2 \rangle} = \sqrt{\frac{3k_B T}{m}}$$
$$m = \frac{0.004 \text{ kg mol}^{-1}}{N_A} \approx 6.6 \times 10^{-27} \text{ kg}$$

$$v_{rms} = \sqrt{\frac{3 \times 1.38 \cdot 10^{-23} \text{ J/K} \times 54.4 \text{ K}}{6.6 \times 10^{-27} \text{ kg}}} \approx 584 \text{ ms}^{-1}$$

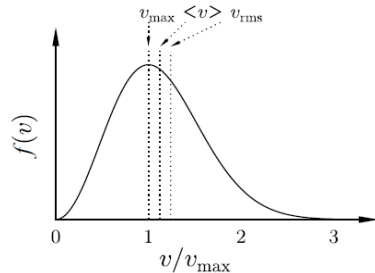
c) The internal energy of a (monoatomic) ideal gas is sum of the particle kinetic energies.

$$\langle u \rangle = \frac{1}{2} m \langle v^2 \rangle$$
$$\langle U \rangle = N \langle u \rangle = N \frac{1}{2} m \langle v^2 \rangle = N \frac{1}{2} m \frac{3k_B T}{m} = \frac{3}{2} N k_B T \approx 225 \text{ J}$$

d)

The temperature would not change, as the number of particles stays the same. The rms speed, however, does depend on $\frac{1}{\sqrt{m}}$. The mean kinetic energy is defined as $\langle E \rangle = \frac{1}{2} m v_{rms}^2 = \frac{3}{2} k_B T$ and does not depend on m , so neither does U .

e)



$$f(v) = \frac{4}{\sqrt{\pi}} \left(\frac{m}{2k_B T} \right)^{3/2} v^2 dv e^{-\frac{mv^2}{2k_B T}}$$

$$m = \frac{0.032 \text{ kg mol}^{-1}}{N_A} \approx 5.3 \times 10^{-26} \text{ kg}$$

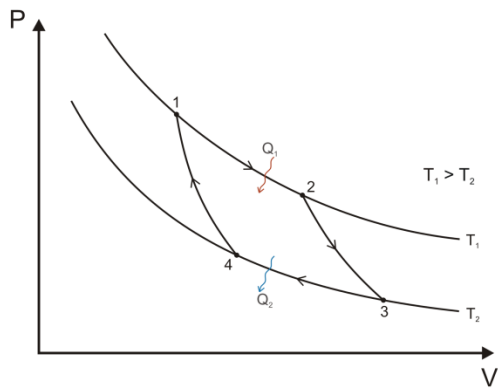
$$f(1550 \text{ ms}^{-1}) = \frac{4}{\sqrt{\pi}} \left(\frac{5.3 \cdot 10^{-26} \text{ kg}}{2k_B 300 \text{ K}} \right)^{3/2} (1050 \text{ ms}^{-1})^2 dv e^{-\frac{5.3 \cdot 10^{-26} \text{ kg} (1050 \text{ ms}^{-1})^2}{2k_B 300 \text{ K}}}$$
$$\approx 1.2 \cdot 10^{-4} \text{ sm}^{-1} dv$$

$$\Delta v = (1500 - 1600) \text{ ms}^{-1} = 100 \text{ ms}^{-1}$$

The fraction is approximately 0.012%.

PROBLEM 2

a)



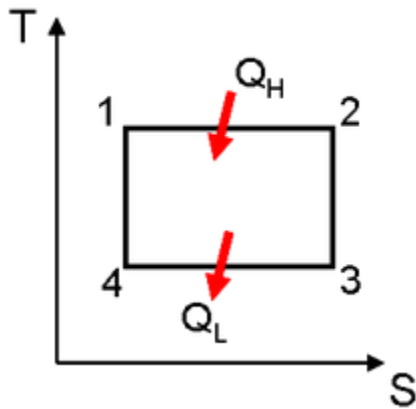
Heat into the engine: i)

Heat out of engine: iii)

Work by engine: i), ii)

Work on engine: iii) iv)

b)



Entropy changes in the isothermal steps i) and iii) (heat flows to keep T constant).

Adiabatic means no heat flow, i.e. no entropy change.

$$\Delta S_A = \int \frac{dQ}{T_h} = \frac{\Delta Q}{T_h} = \frac{Q_h}{T_h}$$

$$\Delta S_B = \frac{Q_c}{T_c}$$

c)

$$\Delta S_h = \Delta S_c$$

$$\frac{Q_h}{T_h} = \frac{Q_c}{T_c}$$

$$\eta = \frac{W}{Q_h} = \frac{Q_h - Q_c}{Q_h} = 1 - \frac{Q_c}{Q_h} = 1 - \frac{T_c}{T_h}$$

d)

$$\eta = \frac{12 \text{ TWh}}{3 \cdot 10^9 \text{ kg} \times 9 \text{ kWh kg}^{-1}} \approx 0.44$$

The cooling water is liquid water from outside the plant (e.g. the sea), let's say $T_c = 300 \text{ K}$.

$$\eta = 1 - \frac{T_c}{T_h}$$

$$T_h \approx \frac{T_c}{0.56} = 535 \text{ K}$$

This is for an idealized Carnot process. The actual temperatures are much higher!

Steam engine: $T_h=100^\circ\text{C}=373\text{K}$, $T_c=50^\circ\text{C}=323\text{K}$, $\eta = 1 - \frac{T_c}{T_h} \approx 0.1$

e) The reservoirs will converge to equal temperatures. At the same time, the efficiency will gradually drop to zero.

PROBLEM 3

a)

$$D \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial T}{\partial r} \right) = 0$$

$$r^2 \frac{\partial T}{\partial r} = \text{const}$$

$$T(r) = A + \frac{B}{r}$$

b)

$$T(r_0) = T_{\text{body}} = A + \frac{B}{r_0}$$

$$T(r \rightarrow \infty) = T_{\text{outside}}$$

$$A = T_{\text{outside}}$$

$$B = (T_{\text{body}} - T_{\text{outside}})r_0$$

c)

$$J_r = -\kappa \frac{\partial T}{\partial r} = \kappa r_0 (T_{\text{body}} - T_{\text{outside}}) \frac{1}{r^2}$$

$$J_{r_0} = \kappa r_0 (T_{\text{body}} - T_{\text{outside}}) \frac{1}{r_0^2}$$

$$J_{r_0} 4\pi r_0^2 = 4\pi \kappa r_0 (T_{\text{body}} - T_{\text{outside}}) \frac{1}{r_0^2} \text{ per second}$$

PROBLEM 4

a)

$$\lambda = \frac{1}{\sqrt{2}n\sigma}$$

$$n = \frac{N}{V} = \frac{p}{k_B T} = \frac{1 \times 10^{-5} \text{ Pa}}{1.38 \times \frac{10^{-23} \text{ J}}{\text{K}} \times 298 \text{ K}} \approx 2.4 \times 10^{15} \text{ m}^{-3}$$
$$\sigma = \pi \times (395 \times 10^{-12})^2 = 4.9 \times 10^{-19} \text{ m}^2$$
$$\lambda = \frac{1}{\sqrt{2} \times 2.4 \times 10^{15} \text{ m}^{-3} \times 4.9 \times 10^{-19} \text{ m}^2} = 601 \text{ m}$$

This is huge, because of the low pressure, molecules essentially do not collide with each other in a normal-sized container.

b) Despite the macroscopic diameter of the pump, we are dealing with effusion rather than with a regular Maxwell-Boltzmann distribution, simply because the mean free path is large as compared to the diameter of the opening.